

QFlux: An Open-Source Python Package for Quantum Dynamics Simulations

QFlux: Beyond Lindblad: Non-Markovian Open Quantum Systems

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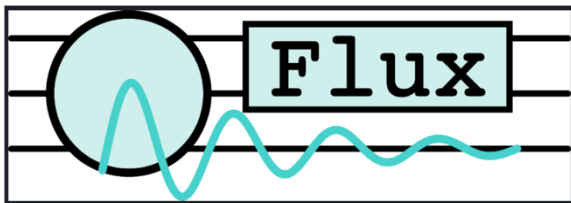
Part VI: Generalized Quantum Master Equation (GQME)

Non-Markovian dynamics, where the environment has memory

Memory kernels + quantum dilation

<https://qflux.batistalab.com>

[JCTC VI.ipynb](#)



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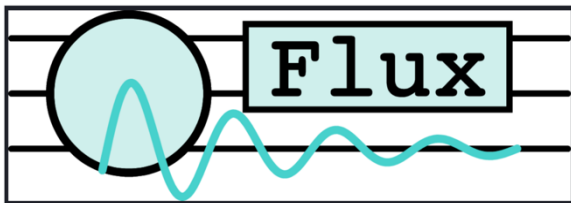
This tutorial is based on the manuscript

QFlux: Quantum Circuit Implementations for Molecular Dynamics

Part VI - The Generalized Quantum Master Equation

Authors:

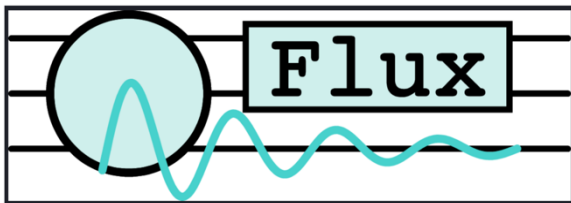
Xiaohan Dan, Pouya Khazaei, Ningyi Lyu, Callie Wilson, Ellen Mulvihill, Yuchen Wang, Brandon C. Allen, Delmar G. A. Cabral, Saurabh Shivpuje, Sabre Kais, Victor S. Batista, and Eitan Geva



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Why Markovian Dynamics Fails?

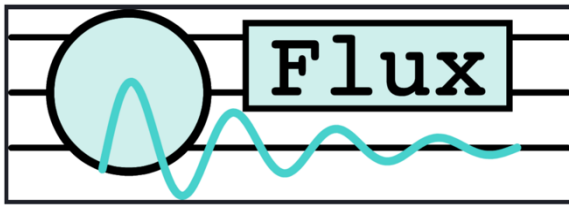
- Lindblad assumes memoryless bath
- Real environments have finite correlation times
- Examples:
 - Solvent relaxation
 - Vibronic coupling
 - Long-lived coherences



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What Is Non-Markovian Dynamics?

- Future depends on the entire past
- Integro-differential equation
- Environment feeds information back



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The GQME: Exact Reduced Dynamics

$$\frac{\partial \hat{\rho}(t)}{\partial t} = -\frac{i}{\hbar} [\hat{H}, \hat{\rho}(t)] = -\frac{i}{\hbar} \mathcal{L} \hat{\rho}(t)$$

$$\hat{H} = \sum_{j=1}^{\infty} \hat{H}_j(\hat{\mathbf{R}}, \hat{\mathbf{P}}) |j\rangle\langle j| + \sum_{j,k=1}^{N_e} \hat{V}_{jk}(\hat{\mathbf{R}}) |j\rangle\langle k|$$

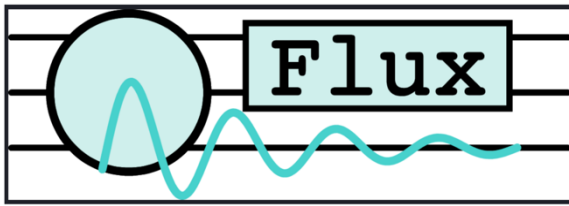
$$\hat{H}_j(\hat{\mathbf{R}}, \hat{\mathbf{P}}) = \hat{\mathbf{P}}^2/2 + V_j(\hat{\mathbf{R}})$$

$$\mathcal{P} \hat{A} = \hat{\rho}_n(0) \otimes \text{Tr}_n\{\hat{A}\}, \quad \mathcal{Q} = \mathcal{I} - \mathcal{P}$$

- GQME:**

$$\dot{\sigma}(t) = \underbrace{-\frac{i}{\hbar} \langle L \rangle_n^0 \sigma(t)}_{\text{Instantaneous bath MF}} - \underbrace{\int_0^t K(\tau) \sigma(t-\tau) d\tau}_{\text{memory term bath influence}}$$

$$\langle \mathcal{L} \rangle_n^0(\cdot) = \text{Tr}_n\{\hat{\rho}_n(0) \mathcal{L}\}(\cdot) \quad \mathcal{K}(\tau) = \frac{1}{\hbar^2} \text{Tr}_n\{\mathcal{L} e^{-i\mathcal{Q}\mathcal{L}\tau/\hbar} \mathcal{Q}\mathcal{L} \hat{\rho}_n(0)\}$$

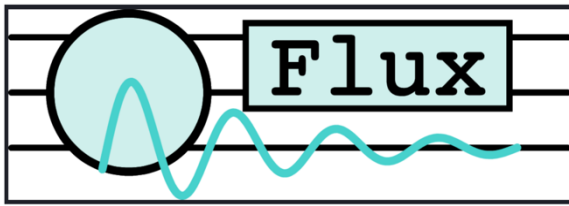


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Strategy of GQME Simulations

Complete workflow

1. From Hamiltonian of microscopic model (e.g., spin–boson)
2. Compute exact reference (TT-TFD)
3. Build projected Liouvillian
4. Compute memory kernel
5. Solve GQME classically
6. Embed non-unitary propagator into a unitary circuit



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Benchmark System: The Spin–Boson Model

$$\hat{H} = \epsilon \hat{\sigma}_z + \Gamma \hat{\sigma}_x + \sum_{i=1}^{N_n} \left[\frac{\hat{P}_i^2}{2} + \frac{1}{2} \omega_i^2 \hat{R}_i^2 - c_i \hat{R}_i \hat{\sigma}_z \right]$$

- Two-level system + harmonic bath.

- **System Parameters:** [JCTC VI.ipynb](#)

Script S.1.1

- bias ϵ
- coupling Γ
- temperature β

- spectral density $J(\omega) = \frac{\pi}{2} \sum_{k=1}^{N_n} \frac{c_k^2}{\omega_k} \delta(\omega - \omega_k)$

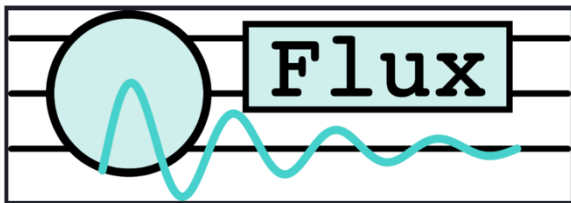
Initial thermal state:

$$\hat{\rho}(0) = |D\rangle\langle D| \otimes \hat{\rho}_n(0),$$

$$\hat{\rho}_n(0) = \frac{\exp \left[-\beta \sum_{i=1}^{N_n} \left(\frac{\hat{P}_i^2}{2} + \frac{1}{2} \omega_i^2 \hat{R}_i^2 \right) \right]}{\text{Tr}_n \left\{ \exp \left[-\beta \sum_{i=1}^{N_n} \left(\frac{\hat{P}_i^2}{2} + \frac{1}{2} \omega_i^2 \hat{R}_i^2 \right) \right] \right\}}$$

Ohmic spectral density:

$$J(\omega) = \frac{\pi \hbar}{2} \xi \omega e^{-\omega/\omega_c}$$



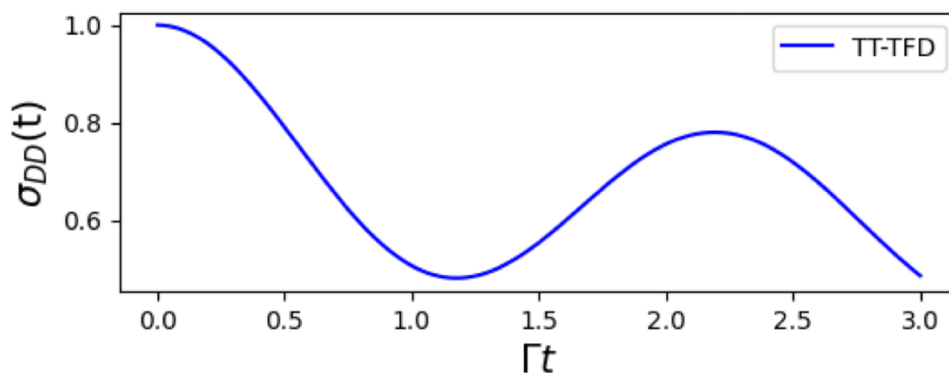
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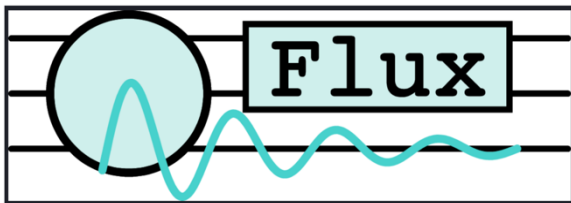
[JCTC VI.ipynb](#)

[Script S.1.2](#)

Reference Dynamics: TT-TFD

- Numerically exact tensor-train method
- Produces $\sigma(t)$ benchmark
- Shows oscillations + dissipation





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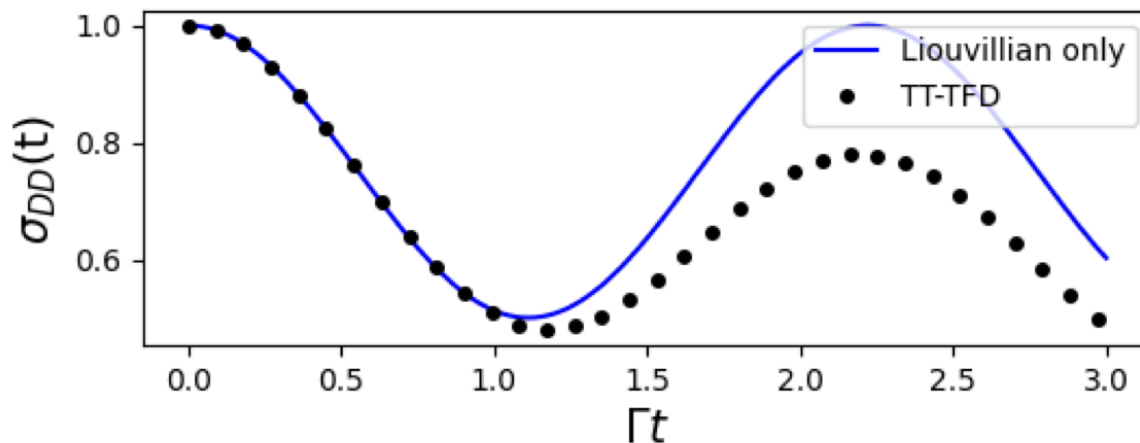
Step 1: Projected Liouvillian (No Memory)

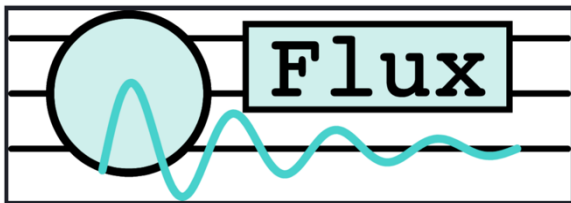
[JCTC VI.ipynb](#)

[Script S.2.1](#)

[Script S.2.1](#)

- Mean-field generator $\langle L \rangle_n^0$
- Equivalent to closed-system dynamics
- Produces undamped oscillations





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Step 2: Building the Memory Kernel

- Solve Volterra equation : [JCTC VI.ipynb](#) [Script S.3.4](#)

$$\mathcal{K}(t) = i\dot{\mathcal{F}}(t) - \frac{1}{\hbar}\mathcal{F}(t)\langle\mathcal{L}\rangle_n^0 + i\int_0^t d\tau \mathcal{F}(t-\tau)\mathcal{K}(\tau)$$

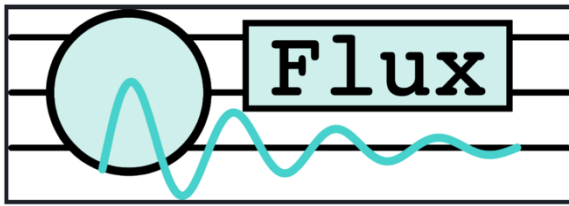
- Projection Free Inputs obtained by TT-TFD: [Script S.3.3](#) [JCTC VI.ipynb](#)

$$\mathcal{F}(t) = \frac{1}{\hbar} \text{Tr}_n \left[\mathcal{L} e^{-i\mathcal{L}t/\hbar} \hat{\rho}_n(0) \right] \quad \dot{\mathcal{F}}(t) = -\frac{i}{\hbar^2} \text{Tr}_n \left[\mathcal{L} e^{-i\mathcal{L}t/\hbar} \mathcal{L} \hat{\rho}_n(0) \right]$$

$$\hat{\sigma}(t) = \underbrace{\mathcal{G}(t)} \hat{\sigma}(0) = \text{Tr}_n \left[e^{-i\mathcal{L}t/\hbar} \hat{\rho}_n(0) \right] \hat{\sigma}(0)$$

[Script S.3.2](#)

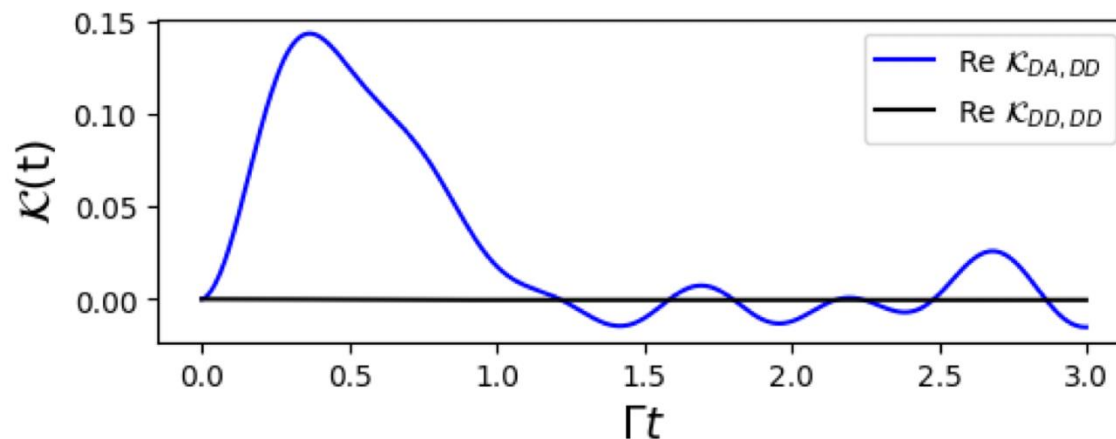
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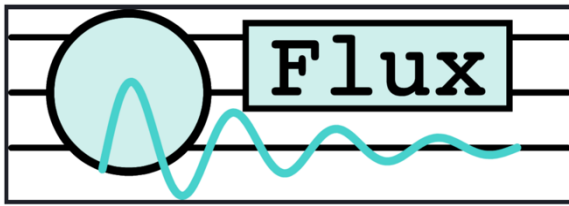


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Interpreting the Kernel

- Some elements small $\rightarrow$ weak direct damping
- Others large $\rightarrow$ population–coherence transfer
- Kernel decays with bath memory time τ_B

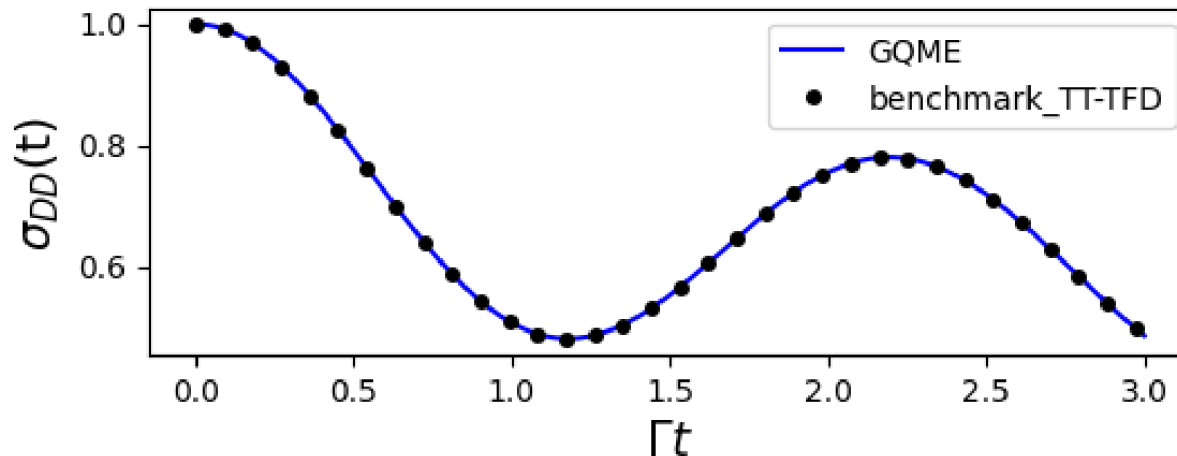


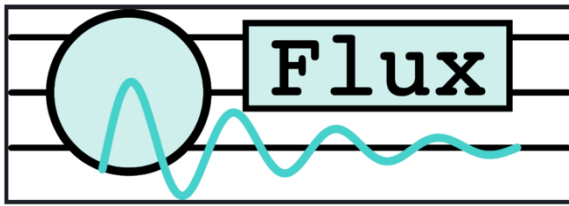


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Step 3: Baseline for the GQME

- RK4 time integration: [JCTC VI.ipynb](#) [Script S.4.2](#)
- Matches TT-TFD exactly





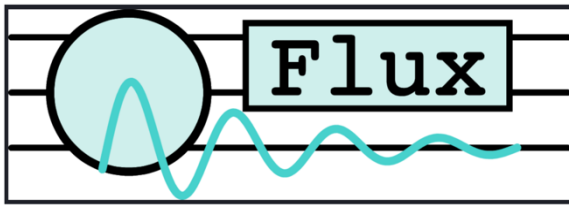
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Problem: GQME Is Non-Unitary

- Obtain propagator by TT-TFD:

$$\frac{\partial \mathcal{G}(t)}{\partial t} = -\frac{i}{\hbar} \langle \mathcal{L} \rangle_n^0 \mathcal{G}(t) - \int_0^t d\tau \mathcal{K}(\tau) \mathcal{G}(t - \tau)$$

- Propagation: $\hat{\sigma}(t) = \mathcal{G}(t) \hat{\sigma}(0)$
- Propagator $\mathcal{G}(t)$ is not unitary
- Quantum circuits require unitary gates
- Need embedding

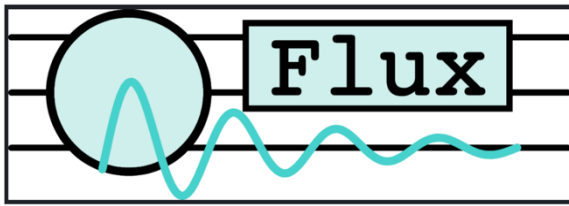


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Key Idea: Sz.-Nagy Dilation

$$\mathcal{U}_{\mathcal{G}'}(t) = \begin{pmatrix} \mathcal{G}'(t) & \mathcal{D}_{\mathcal{G}'^\dagger}(t) \\ \mathcal{D}_{\mathcal{G}'}(t) & -\mathcal{G}'^\dagger(t) \end{pmatrix} \quad \begin{aligned} \mathcal{D}_{\mathcal{G}'}(t) &= \sqrt{I - \mathcal{G}'^\dagger(t)\mathcal{G}'(t)} \\ \mathcal{D}_{\mathcal{G}'^\dagger}(t) &= \sqrt{I - \mathcal{G}'(t)\mathcal{G}'^\dagger(t)} \end{aligned}$$

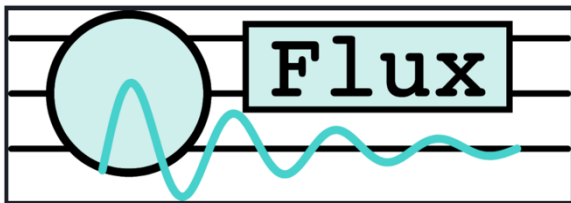
$$\hat{\sigma}(t) = \mathcal{G}'(t)\hat{\sigma}(0) \xrightarrow{\text{dilation}} \mathcal{U}_{\mathcal{G}'}(t) \begin{pmatrix} \hat{\sigma}(0)^\top \\ 0 \\ \vdots \\ 0 \end{pmatrix}$$



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Quantum Workflow

1. Solve GQME $\rightarrow$ get $\mathcal{G}(t)$
2. Rescale to contraction
3. Build dilated unitary $\mathcal{U}_{\mathcal{G}'}(t)$
4. Compile into quantum gate
5. Measure populations



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Circuit Architecture

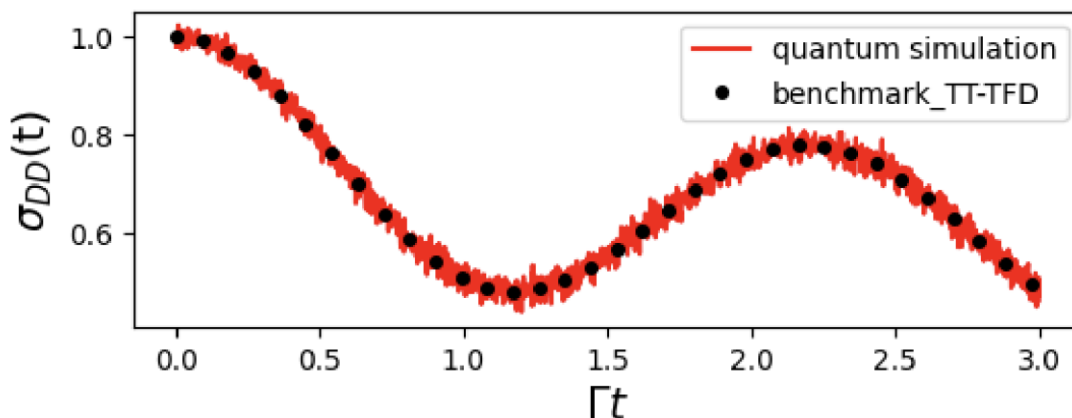
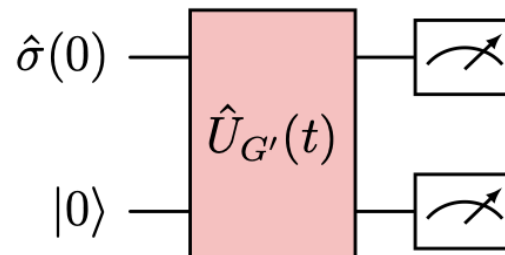
Compact Circuit

[JCTC VI.ipynb](#)

[Script S.6.1](#)

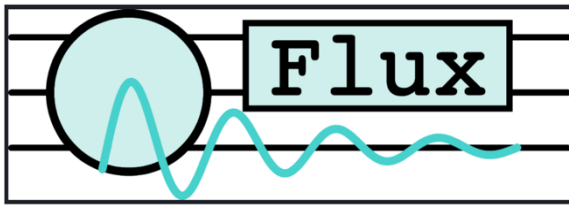
[Script S.6.2](#)

- 3 qubits total
 - 2 system (vectorized density matrix)
 - 1 dilation ancilla
 - One global unitary gate



[JCTC VI.ipynb](#)

[Script S.6.3](#)

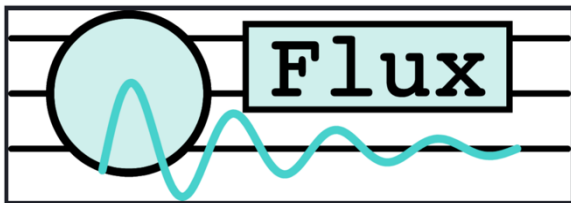


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Quantum Results vs Exact

Memory Effects: Non-Markovian Quantum Simulations

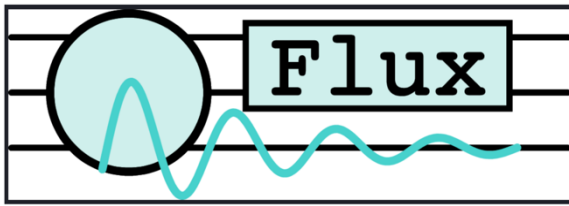
- QASM simulation
- Excellent agreement with TT-TFD
- Noise only from sampling



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Big Picture Take-Home

- GQME gives formally exact reduced dynamics
- Memory kernels capture bath history
- Sz.-Nagy dilation enables quantum simulation
- QFlux provides an end-to-end workflow

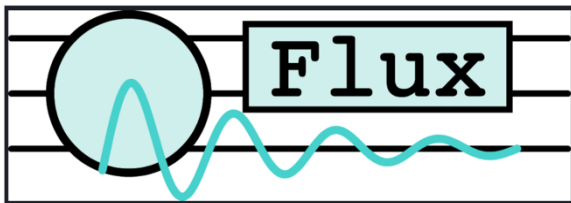


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Conceptual Unification

All simulation methods unified by Qflux under unitary embeddings

- Closed systems $\rightarrow$ unitary evolution
- Markov open systems $\rightarrow$ Lindblad + dilation
- Non-Markov $\rightarrow$ GQME + memory + dilation



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Closing the Series (for now)

- Part I: Classical and Quantum Workflows
- Part II–III: Closed systems
- Part IV–V: Markov open systems
- Part V: Variational methods
- Part VI: non-Markovian memory